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On the Minkowski Measurability of Self-Similar Fractals in \mathbb{R}^d

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Abstract

M. Lapidus and C. Pomerance (1990-1993) and K.J. Falconer (1995) proved that a self-similar fractal in \mathbb{R} is Minkowski-measurable iff it is of non-lattice type. D. Gatzouras (1999) proved that a self-similar fractal in \mathbb{R}^d is Minkowski measurable if it is of non-lattice type (though the actual computation of the content is intractable with his approach) and conjectured that it is not Minkowski measurable if it is of lattice type. Under mild conditions we prove this conjecture and in the non-lattice case we improve his result in the sense that we express the content of the fractal in terms of the residue of the associated ζ -function at the Minkowski-dimension.

Keywords: Self-similar fractals, Minkowski measurability, tube formulas

1 Introduction

Let

$$F = \bigcup_{j=1}^J \varphi_j(F) =: \Phi(F) \subset \mathbb{R}^d$$

be a self-similar fractal, where $\varphi_j : \mathbb{R}^d \rightarrow \mathbb{R}^d$ are similitudes with scaling ratios $0 < r_j < 1$, $j = 1, 2, \dots, J$, for $J \geq 2$. We assume the iterated function system (IFS) Φ to satisfy the open set condition, so that the Minkowski dimension D of F is given by the unique real root of the Moran equation $\sum_{j=1}^J r_j^D = 1$.

Let $F_\varepsilon = \{x \in \mathbb{R}^d \mid \text{dist}(x, F) \leq \varepsilon\}$ and $V_F(\varepsilon)$ be the d -dimensional volume of F_ε . F is called Minkowski measurable if the limit

$$\mathcal{M}(F) := \lim_{\varepsilon \rightarrow 0^+} V_F(\varepsilon) \varepsilon^{D-d}$$

exists, is finite and different from zero. $\mathcal{M}(F)$ is then called the Minkowski content of F .

The IFS Φ is called of lattice type, if the additive subgroup $\sum_{j=1}^J (\log r_j) \mathbb{Z}$ of \mathbb{R} is discrete and otherwise (i.e. if this subgroup is dense in \mathbb{R}) of non-lattice type (see [5]). In the lattice case there is an r with $\log r_j = k_j \log r$, $k_j \in \mathbb{Z}^+$. This dichotomy is decisive for Minkowski measurability of fractals as shown for the one-dimensional case by Lapidus-Pomerance [7],[8] and Falconer [3]. We now briefly recall their results:

Let $d = 1$ and I denote the convex hull of F , $I = [F]$. By the open set condition, $\varphi_j(I)$ and $\varphi_k(I)$ are disjoint for $j \neq k$, except possibly at the endpoints. There will emerge $Q \leq J - 1$ gaps

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on I , with lengths l_q , $q = 1, 2, \dots, Q$. Then, F is Minkowski measurable if and only if the IFS Φ is of non-lattice type, and in case it is measurable the content is given by

$$\mathcal{M} = \frac{2^{1-D} \sum_{q=1}^Q l_q^D}{D(1-D) \sum_{j=1}^J r_j^D \log r_j^{-1}} \quad (\text{see [5, p.262]}). \quad (1)$$

In the lattice case the fractal is not Minkowski measurable, but one can define an average Minkowski content by the formula

$$\mathcal{M}_{\text{av}} = \lim_{T \rightarrow \infty} \frac{1}{\log T} \int_{1/T}^1 \varepsilon^{-(1-D)} V_F(\varepsilon) \frac{d\varepsilon}{\varepsilon} \quad (\text{see [5, p.257]}) \quad (2)$$

and the formula (1) gives in this case the average Minkowski content. (In \mathbb{R}^d , the exponent $-(1-D)$ of ε should be replaced by $-(d-D)$). In higher dimensions, Gatzouras proved in [4] that a non-

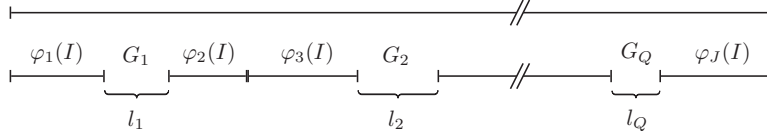


Figure 1:

lattice self-similar fractal in \mathbb{R}^d (satisfying the open set condition) is Minkowski measurable. He used renewal theory and the formula he gave for the Minkowski content is hardly usable for explicit computations. For lattice case he conjectured that the Minkowski content does not exist. Lapidus and van Frankenhuijsen [5, Remark 12.19] remarked that renewal theory was unlikely to yield this result but their approach by higher dimensional tube formulas would settle this issue. Our aim in this paper is in a certain sense to carry out this program under some mild additional conditions on the IFS we give below. Using the Lapidus-Pearse theory [6] we will give an alternative proof of the existence of the Minkowski content in non-lattice case with a very explicit and applicable formula for the content and we will prove the non-Minkowski measurability in the lattice case.

Let $C = [F]$ be the convex hull of F , for which we assume $\dim C = d$. Adopting the approach of Pearse-Winter [10] we want to put some additional conditions on the IFS Φ :

TSC (Tileset Condition): Φ satisfies the open set condition with $\text{int } C$ as a feasible open set.

NTC (Nontriviality Condition): $\text{int } C \not\subseteq \Phi(C) = \bigcup_{j=1}^J \varphi_j(C)$.

Now define $T_0 = C \setminus \Phi(C)$ and its iterates $T_n = \Phi^n(T_0)$, $n = 1, 2, 3, \dots$ (see [9]). The tiling of the self-similar system is given by

$$\mathcal{T} := \{T_n\}_{n=0}^{\infty}.$$

Let $V_{T_n}^-(\varepsilon)$ denote the volume of the inner ε -neighborhood of T_n (i.e. $\{x \in \overline{T_n} \mid \text{dist}(x, T_n^c) \leq \varepsilon\}$) and $V_{\mathcal{T}}^-(\varepsilon) := \sum_{n=0}^{\infty} V_{T_n}^-(\varepsilon)$.

Pearse and Winter prove in [10] the following implication: If the above conditions TSC and NTC hold, then the property $\partial C \subset F$ implies $V_F(\varepsilon) = V_{\mathcal{T}}^-(\varepsilon) + V_C(\varepsilon) - V_C(0)$. This is extremely important, because there are formulas available for $V_{\mathcal{T}}^-(\varepsilon)$ (see below, Theorem 1) and this relationship enables one to compute the true volume of the ε -neighborhood of the fractal. We will call this condition the Pearse-Winter condition:

PWC (Pearse-Winter Condition): $\partial C \subset F$.

To state the tube formula we need some additional assumptions and definitions. Assume that T_0 is the union of finitely many (connected) components, $T_0 = G_1 \cup G_2 \cup \dots \cup G_Q$, called the generators of the tiling. We assume the generators to be (diphase) Steiner-like in the following sense of Lapidus-Pearse [6]: A bounded, open set $G \subset \mathbb{R}^d$ is called (diphase) Steiner-like if the volume $V_G^-(\varepsilon)$ of the inner ε -neighborhood of G admits an expression of the form

$$V_G^-(\varepsilon) = \sum_{m=0}^{d-1} \kappa_m(G) \varepsilon^{d-m}, \quad \text{for } \varepsilon < g, \quad (3)$$

where g denotes the inradius of G , i.e. supremum of the radii of the balls contained in G . For $\varepsilon \geq g$ we have $V_G^-(\varepsilon) = \text{volume}(G)$ which is denoted by $-\kappa_d(G)$, the negative sign being conventional [6].

Lapidus-Pearse introduce the following “scaling ζ -function”:

Definition 1 *The scaling ζ -function of the self-similar fractal is defined by*

$$\zeta(s) = \sum_{n=0}^{\infty} \sum_{w \in W_n} r_w^s,$$

where W_n is the set of words $w = w_1 w_2 \dots w_n$ of length n (with letters from $\{1, 2, \dots, J\}$) and $r_w = r_{w_1} r_{w_2} \dots r_{w_n}$.

The above series can be shown to converge for $\text{Re}(s) > D$. A simple calculation shows that $\zeta(s)$ can be expressed as [5, Theorem 2.4]

$$\zeta(s) = \frac{1}{1 - \sum_{j=1}^J r_j^s} \quad \text{for } \text{Re}(s) > D.$$

$\zeta(s)$ can then be meromorphically extended to the whole complex plane. We will denote this extension also by $\zeta(s)$.

Definition 2 *The set $\mathfrak{D} := \{\omega \in \mathbb{C} \mid \zeta(s) \text{ has a pole at } \omega\}$ is called the set of complex dimensions of the self-similar fractal.*

Lapidus-Pearse define a second type of “ ζ -functions” associated with the tiling and related to the geometry of the (diphase) Steiner-like generators. We assume for simplicity that there is a single generator G (so that $T_0 = G$).

Definition 3 *The geometric ζ -function $\zeta_{\mathcal{T}}(s, \varepsilon)$ associated with the generator G is defined by*

$$\zeta_{\mathcal{T}}(s, \varepsilon) := \zeta(s) \varepsilon^{d-s} \sum_{m=0}^d \frac{g^{s-m}}{s-m} \kappa_m(G).$$

We now state the formula of Lapidus-Pearse for $V_{\mathcal{T}}^-(\varepsilon)$ (though we will use a modification of this formula in the proof of our higher dimensional content formulas below):

Theorem 1 (Tube formula for tilings of self-similar fractals, [6])

$$V_{\mathcal{T}}^-(\varepsilon) = \sum_{\omega \in \mathfrak{D}_{\mathcal{T}}} \text{res}(\zeta_{\mathcal{T}}(s, \varepsilon); \omega),$$

where $\mathfrak{D}_{\mathcal{T}} = \mathfrak{D} \cup \{0, 1, \dots, d-1\}$.

Remark 1 *Lapidus-Pearse gives in [6] a distributional proof for this formula. For a pointwise proof see [1].*

2 Main Results

Our main result is the following theorem:

Theorem 2 Let $F = \Phi(F) = \bigcup_{j=1}^J \varphi_j(F)$ be a self-similar fractal in \mathbb{R}^d with $\dim[F] = d$ and the contractivity ratios of the similitudes $\{\varphi_j\}$ being $\{r_j\}$.

We assume the Tileset Condition, the Nontriviality Condition and the Pearse-Winter Condition to hold (see TSC, NTC and PWC in the former section). We additionally assume $D > d - 1$, where D is the Minkowski dimension of F .

Under these assumptions the followings hold:

I. If the IFS Φ is of non-lattice type, then F is Minkowski measurable with Minkowski content

$$\begin{aligned} \mathcal{M}(F) = \text{res}(\zeta_{\mathcal{T}}(s, \varepsilon) \varepsilon^{s-d}; D) &= \text{res}(\zeta(s); D) \sum_{m=0}^d \frac{g^{D-m}}{D-m} \kappa_m(G) \\ &= \frac{\sum_{m=0}^d \frac{g^{D-m}}{D-m} \kappa_m(G)}{\sum_{j=1}^J r_j^D \log r_j^{-1}}, \end{aligned}$$

II. If the IFS Φ is of lattice type then F is not Minkowski measurable. The average Minkowski content as defined in (2) exists and equals

$$\mathcal{M}_{\text{av}}(F) = \text{res}(\zeta_{\mathcal{T}}(s, \varepsilon) \varepsilon^{s-d}; D).$$

Remark 2 In case of multiple generators, let $\zeta_{\mathcal{T}}^q(s, \varepsilon)$ ($q = 1, 2, \dots, Q$) denote the geometric zeta function corresponding to generator G_q and define the total geometric zeta function by

$$\zeta_{\mathcal{T}}(s, \varepsilon) = \sum_{q=1}^Q \zeta_{\mathcal{T}}^q(s, \varepsilon) = \zeta(s) \varepsilon^{d-s} \sum_{q=1}^Q \sum_{m=0}^d \frac{g_q^{s-m}}{s-m} \kappa_m(G_q),$$

where g_q denotes the inradius of G_q . Then

I. If the IFS Φ is of non-lattice type then F is Minkowski measurable with Minkowski content

$$\mathcal{M}(F) = \text{res}(\zeta_{\mathcal{T}}(s, \varepsilon) \varepsilon^{s-d}; D) = \frac{\sum_{q=1}^Q \sum_{m=0}^d \frac{g_q^{D-m}}{D-m} \kappa_m(G_q)}{\sum_{j=1}^J r_j^D \log r_j^{-1}},$$

II. If the IFS Φ is of lattice type then F is not Minkowski measurable and the average Minkowski content exists and equals

$$\mathcal{M}_{\text{av}}(F) = \sum_{q=1}^Q \text{res}(\zeta_{\mathcal{T}}^q(s, \varepsilon) \varepsilon^{s-d}; D).$$

Corollary 1 If we specialize to the dimension $d = 1$, we obtain the formula (1): Since each G_q is an interval of length l_q we have

$$V_{G_q}^-(\varepsilon) = \begin{cases} 2\varepsilon & , \text{ for } \varepsilon < g_q \\ l_q & , \text{ for } \varepsilon \geq g_q, \end{cases}$$

so that $\kappa_0(G_q) = 2$, $\kappa_1(G_q) = -l_q$ and the inradius g_q equals $l_q/2$. Then

$$\zeta_7^q(s, \varepsilon) = \zeta(s) \varepsilon^{1-s} \left(2 \frac{g_q^s}{s} - l_q \frac{g_q^{s-1}}{s-1} \right) = \zeta(s) \varepsilon^{1-s} \frac{2^{1-s} l_q^s}{s(1-s)},$$

and

$$\mathcal{M}(F) = \text{res} \left(\zeta(s) \frac{2^{1-s}}{s(1-s)} \sum_{q=1}^Q l_q^s; D \right) = \frac{2^{1-D} \sum_{q=1}^Q l_q^D}{D(1-D) \sum_{j=1}^J r_j^D \log r_j^{-1}}.$$

Example 1 Let $\triangle ABC$ be an acute triangle with corresponding sides a, b and c . Let $\triangle A'B'C'$ be its orthic (pedal) triangle (see Fig.2a). The triangles $\triangle AC'B'$, $\triangle BA'C'$ and $\triangle CB'A'$ are scaled copies of the original triangle $\triangle ABC$ with scaling ratios $\cos A, \cos B$ and $\cos C$ (denoting the angle at the vertices A, B, C again with the same letter). Consider the collection of these maps as an iterated function system $\Phi = \{\varphi_j\}_{j=1}^3$ on \mathbb{R}^2 as indicated in Fig.2b. The associated self-similar fractal (“orthic fractal”) is shown in Fig.2c (see [12]). The Minkowski dimension D is determined by $(\cos A)^D + (\cos B)^D + (\cos C)^D = 1$. This system satisfies the TSC, NTC and PWC, and has a single generator $G = \triangle A'B'C'$.

The volume of the inner ε -neighborhood of the generator G is given by

$$V_G(\varepsilon) = \begin{cases} \kappa_1(G)\varepsilon + \kappa_0(G)\varepsilon^2 & , \text{ for } \varepsilon \leq g \\ -\kappa_2(G) & , \text{ for } \varepsilon \geq g \end{cases}$$

where

$$g = \frac{4\text{Area}(\triangle ABC) \cos A \cos B \cos C}{a \cos A + b \cos B + c \cos C},$$

and $\kappa_0(G) = -(\tan A + \tan B + \tan C)$, $\kappa_1(G) = a \cos A + b \cos B + c \cos C$, $\kappa_2(G) = -2 \cos A \cos B \cos C \text{Area}(\triangle ABC)$. Depending on the angles A, B, C , the orthic fractal may be of lattice or non-

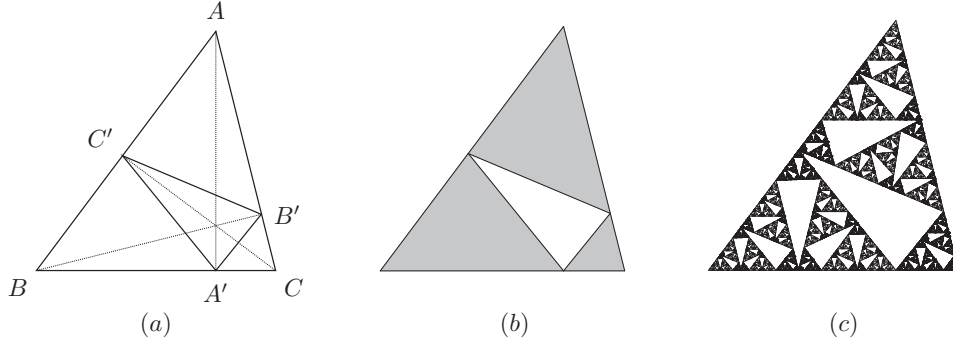


Figure 2: (a) An acute triangle $\triangle ABC$ with its orthic triangle $\triangle A'B'C'$, (b) the IFS, (c) its attractor.

lattice type. If it is of non-lattice type, then by Teorem 2.I, its Minkowski content exists and is given by

$$\frac{\left(\frac{g^D}{D} \kappa_0(G) + \frac{g^{D-1}}{D-1} \kappa_1(G) + \frac{g^{D-2}}{D-2} \kappa_2(G) \right)}{(\cos A)^D \log(\cos A) + (\cos B)^D \log(\cos B) + (\cos C)^D \log(\cos C)}.$$

If the orthic fractal is of lattice type then, by Teorem 2.II, the Minkowski content does not exist, but the average content exists and is given by the same expression.

3 Proof of Theorem 2

We consider first the more difficult non-lattice case (because in that case the distribution of the poles of the ζ -function could be utterly complicated). By the assumptions of the theorem we have

$$V_F(\varepsilon) = V_{\mathcal{T}}^-(\varepsilon) + V_C(\varepsilon) - V_C(0), \text{ where } C = [F].$$

We have to consider the limit behaviour of $V_F(\varepsilon) \varepsilon^{D-d}$ as ε tends to zero.

By the well-known Steiner formula, the volume of the ε -neighborhood of a bounded convex set in \mathbb{R}^d can be expressed as a polynomial in ε [11]:

$$V_C(\varepsilon) = \sum_{m=0}^d a_m \varepsilon^m \quad \text{with } a_0 = V_C(0).$$

Hence, $\lim_{\varepsilon \rightarrow 0^+} (V_C(\varepsilon) - V_C(0)) \varepsilon^{D-d} = 0$ (by the assumption $D > d - 1$). Thus, our concern will be the term $V_{\mathcal{T}}^-(\varepsilon) \varepsilon^{D-d}$.

To make the proof transparent, we will formulate several lemmas, whose proofs we defer to the next section.

Lemma 1 *For any c satisfying $D < c < d$,*

$$V_{\mathcal{T}}^-(\varepsilon) = \frac{1}{2\pi\mathbf{i}} \int_{c-\mathbf{i}\infty}^{c+\mathbf{i}\infty} \zeta_{\mathcal{T}}(s, \varepsilon) ds.$$

Now we want to convert this integral into an appropriate sum of residues of $\zeta_{\mathcal{T}}(s, \varepsilon)$ plus an integral on a path Γ lying to the left of the line $\text{Re}(s) = D$.

For the construction of this path Γ we need the following two lemmas. For convenience we assume that the contractivity ratios are ordered as

$$1 > r_1 \geq r_2 \geq \dots \geq r_J > 0.$$

Lemma 2 *There exists $\tilde{D} < D$ such that all the poles of $\zeta(s)$ in the strip $\{s \mid \tilde{D} < \text{Re}(s) < D\}$ are simple and the residues of $\zeta(s)$ at these poles are bounded by $1/\log r_1^{-1}$.*

Lemma 3 *There exist strictly increasing, real sequences $\{\alpha_k\}_{k \in \mathbb{Z}}$ and $\{\beta_k\}_{k \in \mathbb{Z}}$ with $\alpha_k < \beta_k < \alpha_{k+1}$ for all k , $\alpha_0 < 0 < \beta_0$ and*

$$\alpha_{k+1} - \alpha_k > \frac{\pi}{\log r_J^{-1}}, \quad (k \in \mathbb{Z})$$

and there exist σ_L, σ_R with $\max\{\tilde{D}, d-1\} < \sigma_L < \sigma_R < D$, such that $\zeta(s)$ is uniformly bounded for all $k \in \mathbb{Z}$ on the (oriented) segments

$$\begin{aligned} \gamma_k^1 &:= [\sigma_R + \mathbf{i}\beta_{k-1}, \sigma_R + \mathbf{i}\alpha_k], & \gamma_k^2 &:= [\sigma_R + \mathbf{i}\alpha_k, \sigma_L + \mathbf{i}\alpha_k], \\ \gamma_k^3 &:= [\sigma_L + \mathbf{i}\alpha_k, \sigma_L + \mathbf{i}\beta_k], & \gamma_k^4 &:= [\sigma_L + \mathbf{i}\beta_k, \sigma_R + \mathbf{i}\beta_k]. \end{aligned}$$

Let Γ_k be the concatenation of the segments γ_k^l , $l = 1, 2, 3, 4$ and Γ be the path obtained by the concatenation of all Γ_k , $k \in \mathbb{Z}$. Let Ω be the open region between Γ and the line $\text{Re}(s) = D$. Then, by Lemma 2 and Lemma 3, there exists $K > 0$ such that

$$|\zeta(s)| \leq K, \text{ for all } s \in \Gamma$$

and

$$|\text{res}(\zeta(s); \omega)| \leq K \text{ for all poles } \omega \in \Omega \text{ of } \zeta.$$

(As there are too many constants in the sequel, we will use the letter K for any of them, though they may differ in the appearing context.)

As $\zeta(s)$ is analytic in $\{s \mid \text{Re}(s) > D\}$, all the poles of ζ lie in the half plane $\{s \mid \text{Re}(s) \leq D\}$, and by [5, Theorem 2.17], D is the only pole of ζ with real part D . Now, the integral in Lemma 1 can be expressed as follows:

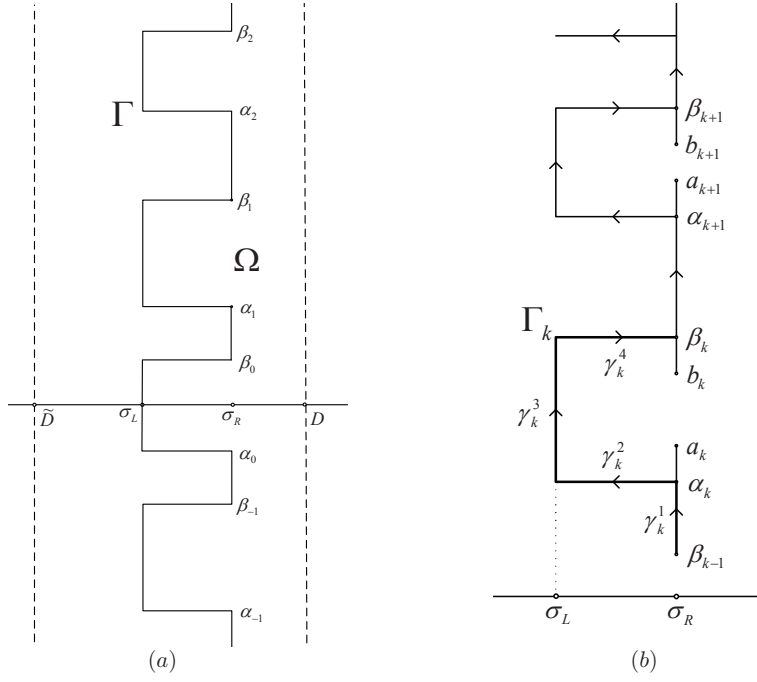


Figure 3: The path Γ

Lemma 4

$$\frac{1}{2\pi i} \int_{c-i\infty}^{c+i\infty} \zeta_{\mathcal{T}}(s, \varepsilon) ds = \text{res}(\zeta_{\mathcal{T}}(s, \varepsilon); D) + \sum_{\omega \in \Omega \cap \mathfrak{D}} \text{res}(\zeta_{\mathcal{T}}(s, \varepsilon); \omega) + \frac{1}{2\pi i} \int_{\Gamma} \zeta_{\mathcal{T}}(s, \varepsilon) ds.$$

The integral over Γ on the right-hand side above is absolutely convergent and can be estimated as follows:

Lemma 5 $\int_{\Gamma} |\zeta_{\mathcal{T}}(s, \varepsilon)| |ds| = O(\varepsilon^{d-\sigma_R})$ as $\varepsilon \rightarrow 0^+$.

This means that we will get rid of this term in the evaluation of the limit $V_{\mathcal{T}}^-(\varepsilon)\varepsilon^{D-d}$ as $\varepsilon \rightarrow 0^+$: $O(\varepsilon^{d-\sigma_R})\varepsilon^{D-d} = o(1)$ since $\sigma_R < D$.

We have

$$\lim_{\varepsilon \rightarrow 0^+} \text{res}(\zeta_{\mathcal{T}}(s, \varepsilon); D)\varepsilon^{D-d} = \frac{\sum_{m=0}^d \frac{q^{D-m}}{D-m} \kappa_m(G)}{\sum_{j=1}^J r_j^D \log r_j^{-1}},$$

where the numerator of the right-hand side is different from zero by Remark 4 below and the denominator is obviously non-zero. Therefore, the proof of first part of Theorem 2 will be settled by the following lemma:

Lemma 6 $\lim_{\varepsilon \rightarrow 0^+} \varepsilon^{D-d} \sum_{\omega \in \Omega \cap \mathfrak{D}} \text{res}(\zeta_{\mathcal{T}}(s, \varepsilon); \omega) = 0.$

Now we consider the lattice case. In this case we can use simply a vertical line $\text{Re}(s) = \sigma < D$ (with σ sufficiently close to D) instead of the complicated Γ and applying the same procedures we

can arrive at the formula

$$V_{\mathcal{T}}^{-}(\varepsilon) = \sum_{n=-\infty}^{\infty} \operatorname{res}(\zeta_{\mathcal{T}}(s, \varepsilon); D + \mathbf{i}np) + \frac{1}{2\pi\mathbf{i}} \int_{\sigma-\mathbf{i}\infty}^{\sigma+\mathbf{i}\infty} \zeta_{\mathcal{T}}(s, \varepsilon) ds,$$

where $p = 2\pi/\log r$ (with $\log r$ being the generator of the group $\sum_{j=1}^J (\log r_j)\mathbb{Z}$). As in Lemma 5, the integral on the right-hand side is $O(\varepsilon^{d-\sigma})$ as $\varepsilon \rightarrow 0^+$, so that we can omit this term. The non-real complex dimensions emerging on the line $\operatorname{Re}(s) = D$ will now cause oscillations and prevent the function $\varepsilon^{D-d}V_{\mathcal{T}}^{-}(\varepsilon)$ to have a limit as $\varepsilon \rightarrow 0^+$:

$$\begin{aligned} \varepsilon^{D-d}V_{\mathcal{T}}^{-}(\varepsilon) &= \frac{1}{\sum_{j=1}^J r_j^D \log r_j^{-1}} \sum_{n \in \mathbb{Z}} \varepsilon^{-\mathbf{i}np} \sum_{m=0}^d \frac{g^{D+\mathbf{i}np-m}}{D + \mathbf{i}np - m} \kappa_m(G) \\ &=: \frac{1}{\sum_{j=1}^J r_j^D \log r_j^{-1}} \sum_{n \in \mathbb{Z}} a_n \varepsilon^{-\mathbf{i}np} = \frac{1}{\sum_{j=1}^J r_j^D \log r_j^{-1}} \sum_{n \in \mathbb{Z}} a_n e^{\mathbf{i}np} \end{aligned}$$

by change of variable $x = -\log \varepsilon$. By 5 below at most $d-1$ of a_n can vanish and by 6 $\sum_{n \in \mathbb{Z}} |a_n| < \infty$, so that the above Fourier series uniformly converges, is non-constant and oscillates as $x \rightarrow \infty$ ($\varepsilon \rightarrow 0^+$). Thus a lattice fractal is never Minkowski measurable. However, the average Minkowski content always exists and can be calculated as follows (as $V_C(\varepsilon) - V_C(0)$ does not contribute):

$$\begin{aligned} \mathcal{M}_{\text{av}} &= \lim_{T \rightarrow \infty} \frac{1}{\log T} \int_{1/T}^1 \varepsilon^{D-d} V_F(\varepsilon) \frac{d\varepsilon}{\varepsilon} = \lim_{T \rightarrow \infty} \frac{1}{\log T} \int_{1/T}^1 \varepsilon^{D-d} V_{\mathcal{T}}^{-}(\varepsilon) \frac{d\varepsilon}{\varepsilon} \\ &= \lim_{T \rightarrow \infty} \frac{1}{\log T} \frac{1}{\sum_{j=1}^J r_j^D \log r_j^{-1}} \sum_{n \in \mathbb{Z}} \int_0^{\log T} a_n e^{\mathbf{i}np} dx = \frac{a_0}{\sum_{j=1}^J r_j^D \log r_j^{-1}}, \end{aligned}$$

where the third equality we use the uniform convergence.

4 Proof of Lemmas

Proof of Lemma 1. Proof of a more general version of this lemma can be found in [2]. For convenience of the reader we repeat the main steps below, omitting the justification of technical details. We have

$$V_{\mathcal{T}}^{-}(\varepsilon) = \sum_{n=0}^{\infty} V_{T_n}^{-}(\varepsilon) = \sum_{n=0}^{\infty} \sum_{w \in W_n} V_{r_w G}^{-}(\varepsilon),$$

where $r_w G$ is a copy of G scaled by r_w . Recall that $V_G^{-}(\varepsilon)$ is given as in 3. A simple calculation shows

$$V_{rG}^{-}(\varepsilon) = \begin{cases} \sum_{m=0}^{d-1} \kappa_m(G) r^m \varepsilon^{d-m} & , \text{ for } \varepsilon < rg \\ -r^d \kappa_d(G) & , \text{ for } \varepsilon \geq rg. \end{cases}$$

We calculate the Mellin transform of $V_{\mathcal{T}}^{-}(\varepsilon)/\varepsilon^d$: The Mellin transform $\mathfrak{M}[f; s]$ of $f : (0, \infty) \rightarrow \mathbb{R}$ is given by $\mathfrak{M}[f; s] = \int_0^{\infty} f(x)x^{s-1} dx$. A routine calculation shows that

$$\mathfrak{M} \left[\frac{V_{rG}^{-}(\varepsilon)}{\varepsilon^d}; s \right] = r^s \sum_{m=0}^d \kappa_m(G) \frac{g^{s-m}}{s-m}, \quad \text{for } d-1 < \operatorname{Re}(s) < d. \quad (4)$$

Therefore, for $D < \operatorname{Re}(s) < d$,

$$\begin{aligned} \mathfrak{M} \left[\frac{V_{\mathcal{T}}^-(\varepsilon)}{\varepsilon^d}; s \right] &= \int_0^\infty \sum_{n=0}^\infty \sum_{w \in W_n} V_{r_w G}^-(\varepsilon) \varepsilon^{s-d-1} ds = \sum_{n=0}^\infty \sum_{w \in W_n} \int_0^\infty V_{r_w G}^-(\varepsilon) \varepsilon^{s-d-1} ds \\ &= \sum_{n=0}^\infty \sum_{w \in W_n} r_w^s \left(\sum_{m=0}^d \kappa_m(G) \frac{g^{s-m}}{s-m} \right) = \zeta(s) \sum_{m=0}^d \kappa_m(G) \frac{g^{s-m}}{s-m}. \end{aligned}$$

Taking the inverse Mellin transform, we obtain

$$\frac{V_{\mathcal{T}}^-(\varepsilon)}{\varepsilon^d} = \frac{1}{2\pi i} \int_{c-i\infty}^{c+i\infty} \left(\zeta(s) \sum_{m=0}^d \kappa_m(G) \frac{g^{s-m}}{s-m} \right) \varepsilon^{-s} ds,$$

where c satisfies $D < c < d$. Hence the claim is proved. ■

In passing we note the following results which will be useful for us later.

Remark 3 Putting $r = 1$ in 4 gives

$$\int_0^\infty V_G^-(\varepsilon) \varepsilon^{s-d-1} d\varepsilon = \sum_{m=0}^d \kappa_m \frac{g^{s-m}}{s-m} = g^{s-d} \sum_{m=0}^d \kappa_m \frac{g^{d-m}}{s-m} = g^{s-d} \frac{P(s)}{Q(s)}, \quad (5)$$

where $Q(s) = s(s-1)\cdots(s-d)$ is a polynomial of degree $d+1$. A crucial observation is that the degree of the polynomial $P(s)$ is at most $d-1$. This is a consequence of the continuity of $V_G^-(\varepsilon)$ at $\varepsilon = g$: $\sum_{m=0}^{d-1} \kappa_m g^{d-m} = -\kappa_d$.

Therefore, for any σ_1, σ_2 with $d-1 < \sigma_1 < \sigma_2 < d$, there exists $K > 0$ such that

$$\left| \sum_{m=0}^d \kappa_m \frac{g^{s-m}}{s-m} \right| \leq \frac{K}{|s|^2} \quad \text{for } \sigma_1 \leq \operatorname{Re}(s) \leq \sigma_2. \quad (6)$$

Remark 4 Putting $s = D$ in 5 gives

$$\sum_{m=0}^d \kappa_m \frac{g^{D-m}}{D-m} = \int_0^\infty V_G^-(\varepsilon) \varepsilon^{D-d-1} d\varepsilon,$$

which shows that $\sum_{m=0}^d \kappa_m \frac{g^{D-m}}{D-m}$ can not be zero.

Proof of Lemma 2. Recall that the contractivity ratios are assumed to be ordered as $1 > r_1 \geq r_2 \geq \cdots \geq r_J > 0$. Let D satisfy the Moran equation: $r_1^D + r_2^D + \cdots + r_J^D = 1$. We define \tilde{D} to be the unique real solution of the equation $r_1^{\tilde{D}} + r_2^{\tilde{D}} + \cdots + r_{J-1}^{\tilde{D}} = 1$. It is clear that $\tilde{D} < D$.

Let $f(s) = 1 - (r_1^s + r_2^s + \cdots + r_J^s)$, so that $\zeta(s) = 1/f(s)$. Let $s_0 = \sigma_0 + it_0$ be a zero of $f(s)$ in the strip $\{s \mid \tilde{D} < \operatorname{Re}(s) < D\}$. We will first show that $\operatorname{Re}(r_j^{s_0}) \geq 0$, for all $j = 1, 2, \dots, J$:

Since s_0 is a zero of f , we have $r_1^{s_0} + r_2^{s_0} + \cdots + r_J^{s_0} = 1$. Taking real parts, we obtain

$$\sum_{j=1}^J \operatorname{Re}(r_j^{s_0}) = 1. \quad (7)$$

Suppose that for some j_0 , we have $\operatorname{Re}(r_{j_0}^{s_0}) < 0$. Then,

$$\sum_{j=1}^J \operatorname{Re}(r_j^{s_0}) < \sum_{j=1, j \neq j_0}^J \operatorname{Re}(r_j^{s_0}) \leq \sum_{j=1, j \neq j_0}^J r_j^{\sigma_0} \leq \sum_{j=1}^{J-1} r_j^{\sigma_0} < \sum_{j=1}^{J-1} r_j^{\tilde{D}} = 1.$$

This contradicts to (7).

The nonnegativity of $\operatorname{Re}(r_j^{s_0})$ and 7 implies that

$$\operatorname{Re}(f'(s_0)) = \sum_{j=1}^J \log r_j^{-1} \operatorname{Re}(r_j^{s_0}) \geq \log r_1^{-1} \sum_{j=1}^J \operatorname{Re}(r_j^{s_0}) = \log r_1^{-1}.$$

Thus, the zero of $f(s)$ (and therefore the pole of $\zeta(s)$) at $s = s_0$ is simple. Moreover,

$$|\operatorname{res}(\zeta(s); s_0)| = \left| \frac{1}{f'(s_0)} \right| \leq \left| \frac{1}{\operatorname{Re}(f'(s_0))} \right| \leq \frac{1}{\log r_1^{-1}}.$$

■ **Proof of Lemma 3.** With \tilde{D} as in Lemma 2, choose $\sigma_L < D$ such that $\sigma_L > \max\{\tilde{D}, d-1\}$. Let $r_1^{\sigma_L} + r_2^{\sigma_L} + \cdots + r_J^{\sigma_L} =: 1 + \lambda$. Let $0 < \psi < \frac{\pi}{2}$ be chosen such that

$$3\psi \frac{|\log r_J|}{|\log r_1|} < \sqrt{\frac{\lambda}{1+\lambda}} \quad (8)$$

and let

$$\mu = \frac{\psi^2 r_J^D}{8}. \quad (9)$$

Then there exists a unique real number σ_R such that $r_1^{\sigma_R} + r_2^{\sigma_R} + \cdots + r_J^{\sigma_R} = 1 + \mu$. Note that

$$\mu < \frac{\psi^2}{8} < 9\psi^2 \frac{|\log r_J|^2}{|\log r_1|^2} < \frac{\lambda}{1+\lambda} < \lambda,$$

hence $\sigma_L < \sigma_R$.

For $s = \sigma + it$, let $-\pi \leq \theta_j(t) < \pi$ be the angle of r_j^s : $\theta_j(t) \equiv t \log r_j \pmod{2\pi}$.

To construct the sequences $\{\alpha_k\}$ and $\{\beta_k\}$, we first determine the points $s = \sigma_R + it$ on the line $\operatorname{Re}(s) = \sigma_R$ for which $|\theta_j(t)| < \psi$ for all $j = 1, 2, \dots, J$.

The set $\{t \mid |\theta_j(t)| < \psi \text{ for all } j = 1, 2, \dots, J\}$ is a union of countably many disjoint open intervals. That is

$$\{t \mid |\theta_j(t)| < \psi \text{ for all } j = 1, 2, \dots, J\} = \bigcup_{k \in \mathbb{Z}} I_k =: \bigcup_{k \in \mathbb{Z}} (a_k, b_k),$$

with $|I_k| = b_k - a_k \leq 2\psi/|\log r_J|$ and $a_{k+1} - b_k \geq (2\pi - 2\psi)/|\log r_J|$.

We define $\alpha_k := a_k - 2\psi/|\log r_1|$ and $\beta_k := b_k + 2\psi/|\log r_1|$ (see Fig.4b). Clearly, $\alpha_k < \beta_k$. The inequality $\beta_k < \alpha_{k+1}$ follows from

$$2 \frac{2\psi}{|\log r_1|} < \frac{2\pi - 2\psi}{|\log r_J|}.$$

This inequality is a consequence of the following inequalities (the second one being 8):

$$\psi \left(\frac{2|\log r_J|}{|\log r_1|} + 1 \right) < 3\psi \frac{|\log r_J|}{|\log r_1|} < \sqrt{\frac{\lambda}{1+\lambda}} < 1 < \pi.$$

Moreover $\alpha_{k+1} - \alpha_k = a_{k+1} - a_k > a_{k+1} - b_k \geq (2\pi - 2\psi)/|\log r_J| > \pi/|\log r_J|$.

We will prove that $\zeta(s)$ is uniformly bounded on the (oriented) segments γ_k^l , ($l = 1, 2, 3, 4$ and $k \in \mathbb{Z}$) (see Fig. 3b). This will follow from the following estimates (recall that $f(s) = 1 - (r_1^s + r_2^s + \cdots + r_J^s)$ and $\zeta(s) = 1/f(s)$):

- i) $\operatorname{Re}(f(s)) \geq \mu$ for $s \in \gamma_k^1$,
- ii) $\operatorname{Im}(f(s)) \leq -\sin \psi$ for $s \in \gamma_k^2$,

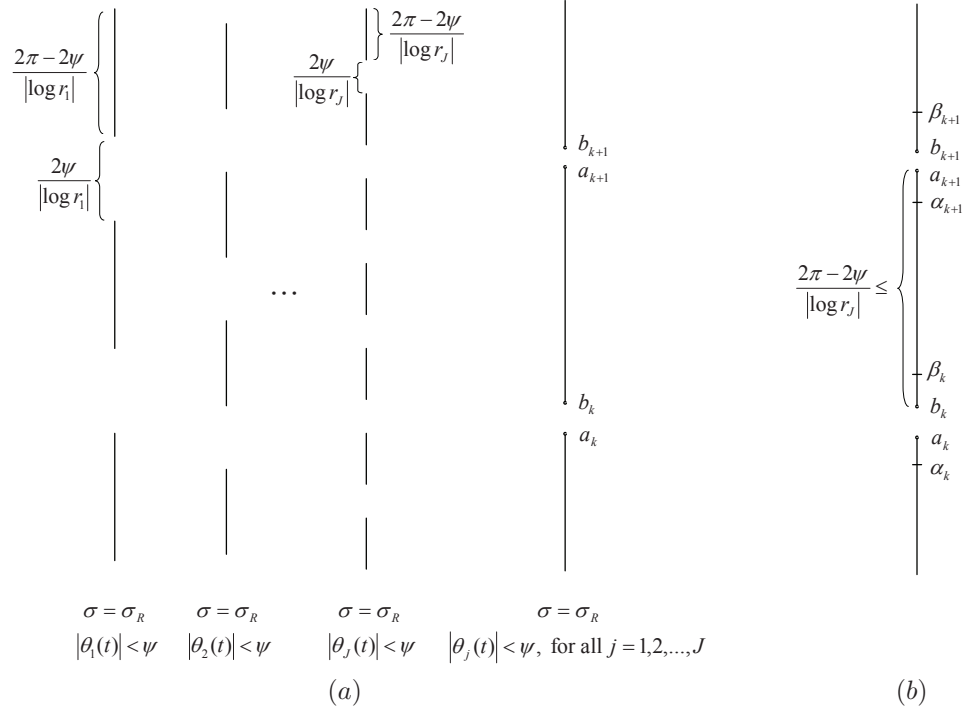


Figure 4: Construction of the sequences $\{a_k\}$, $\{b_k\}$, $\{\alpha_k\}$ and $\{\beta_k\}$.

iii) $\operatorname{Re}(f(s)) \leq -\frac{\lambda}{2}$ for $s \in \gamma_k^3$,

iv) $\operatorname{Im}(f(s)) \geq \sin \psi$ for $s \in \gamma_k^4$.

We begin with i): For $s = \sigma_R + \mathbf{i}t \in \gamma_k^1$, we have $|\theta_j(t)| \geq \psi$ for at least one $j = j_0$. Using 9 and the inequality $\cos \theta \leq 1 - \theta^2/4$ for $-\pi/2 \leq \theta \leq \pi/2$, we get

$$\begin{aligned} \operatorname{Re}(r_{j_0}^s) &= r_{j_0}^{\sigma_R} \cos \theta_{j_0}(t) < r_{j_0}^{\sigma_R} \cos \psi \\ &\leq r_{j_0}^{\sigma_R} \left(1 - \frac{\psi^2}{4}\right) = r_{j_0}^{\sigma_R} \left(1 - \frac{2\mu}{r_j^D}\right) < r_{j_0}^{\sigma_R} - 2\mu \frac{r_{j_0}^D}{r_j^D} \leq r_{j_0}^{\sigma_R} - 2\mu. \end{aligned}$$

Hence,

$$\begin{aligned} \operatorname{Re} \left(\sum_{j=1}^J r_j^s \right) &= \operatorname{Re}(r_{j_0}^s) + \operatorname{Re} \left(\sum_{j=1, j \neq j_0}^J r_j^s \right) \\ &\leq r_{j_0}^{\sigma_R} - 2\mu + \sum_{j=1, j \neq j_0}^J r_j^{\sigma_R} = \sum_{j=1}^J r_j^{\sigma_R} - 2\mu = 1 + \mu - 2\mu = 1 - \mu. \end{aligned}$$

Therefore $\operatorname{Re}(f(s)) \geq \mu$.

We now prove ii): We first show that for each $j = 1, 2, \dots, J$, $\theta_j(\alpha_k)$ satisfies $\psi \leq \theta_j(\alpha_k) \leq \sqrt{\frac{\lambda}{1+\lambda}}$. Fix $j \in \{1, 2, \dots, J\}$. By definition of a_k , we have $-\psi \leq \theta_j(a_k) \leq \psi$. Therefore, there exists $m \in \mathbb{Z}$ such that $2\pi m - \psi \leq a_k \log r_j \leq 2\pi m + \psi$. Then, since $\alpha_k = a_k - 2\psi/|\log r_1|$,

$$2\pi m - \psi + \frac{2\psi}{|\log r_1|} |\log r_j| \leq \alpha_k \log r_j \leq 2\pi m + \psi + \frac{2\psi}{|\log r_1|} |\log r_j|.$$

Using 8 and noting that $|\log r_j| \geq |\log r_1|$, we obtain

$$2\pi m + \psi \leq \alpha_k \log r_j \leq 2\pi m + \sqrt{\frac{\lambda}{1+\lambda}}.$$

Therefore $\psi \leq \theta_j(\alpha_k) \leq \sqrt{\frac{\lambda}{1+\lambda}} < 1 < \frac{\pi}{2}$.

Now, for $s = \sigma + \mathbf{i}\alpha_k \in \gamma_k^2$, we have $\sigma_L \leq \sigma \leq \sigma_R < D$ and

$$\operatorname{Im}(f(s)) = -\sum_{j=1}^J r_j^\sigma \sin \theta_j(\alpha_k) \leq -\sin \psi \sum_{j=1}^J r_j^\sigma \leq -\sin \psi \sum_{j=1}^J r_j^D = -\sin \psi.$$

We now prove iii): Reasoning as we did in the proof of part ii), it can be easily shown that $-\sqrt{\frac{\lambda}{1+\lambda}} \leq \theta_j(\beta_k) \leq -\psi$ and for $\alpha_k \leq t \leq \beta_k$, we have $-\sqrt{\frac{\lambda}{1+\lambda}} \leq \theta_j(t) \leq \sqrt{\frac{\lambda}{1+\lambda}}$, for every $j \in \{1, 2, \dots, J\}$. For $s \in \gamma_k^3$, we have $s = \sigma_L + \mathbf{i}t$, $\alpha_k \leq t \leq \beta_k$. Noting that $\cos \theta \geq 1 - \theta^2/2$,

$$\begin{aligned} \operatorname{Re}(f(s)) &= 1 - \sum_{j=1}^J r_j^{\sigma_L} \cos \theta_j(t) \leq 1 - \cos \left(\sqrt{\frac{\lambda}{1+\lambda}} \right) \sum_{j=1}^J r_j^{\sigma_L} \\ &\leq 1 - \left(1 - \frac{1}{2} \frac{\lambda}{1+\lambda} \right) (1+\lambda) = -\frac{\lambda}{2}. \end{aligned}$$

Finally, we prove the case iv). For $s \in \gamma_k^4$, we have $s = \sigma + \mathbf{i}\beta_k$, $\sigma_L \leq \sigma \leq \sigma_R < D$ and $-\sqrt{\frac{\lambda}{1+\lambda}} \leq \theta_j(\beta_k) \leq -\psi$. Then

$$\operatorname{Im}(f(s)) = -\sum_{j=1}^J r_j^\sigma \sin \theta_j(\beta_k) \geq -\sin(-\psi) \sum_{j=1}^J r_j^\sigma \geq \sin \psi \sum_{j=1}^J r_j^D = \sin \psi.$$

■

Proof of Lemma 4.

By [5, Theorem 3.26], there exists an increasing sequence $\{\rho_n\}_{n=1}^\infty$ tending to infinity such that $\zeta(s)$ is uniformly bounded on the lines $\operatorname{Im}(s) = \pm\rho_n$. That is, there exists $K > 0$ such that

$$|\zeta(s)| \leq K, \quad \operatorname{Im}(s) = \pm\rho_n, \quad n = 1, 2, \dots$$

By residue theorem,

$$\begin{aligned} \frac{1}{2\pi\mathbf{i}} \int_{c-\mathbf{i}\rho_n}^{c+\mathbf{i}\rho_n} \zeta_{\mathcal{T}}(s, \varepsilon) ds &= \frac{1}{2\pi\mathbf{i}} \int_{L_n} \zeta_{\mathcal{T}}(s, \varepsilon) ds + \frac{1}{2\pi\mathbf{i}} \int_{\Gamma_n} \zeta_{\mathcal{T}}(s, \varepsilon) ds \\ &+ \frac{1}{2\pi\mathbf{i}} \int_{L'_n} \zeta_{\mathcal{T}}(s, \varepsilon) ds + \sum_{\omega \in \Omega_{\rho_n} \cap \mathfrak{D}} \operatorname{res}(\zeta_{\mathcal{T}}(s, \varepsilon); \omega) \end{aligned} \quad (10)$$

where $L_n = \{s \mid \operatorname{Im}(s) = \rho_n\} \cap \overline{\Omega}$, $L'_n = \{s \mid \operatorname{Im}(s) = -\rho_n\} \cap \overline{\Omega}$, $\Gamma_n = \{s \mid -\rho_n \leq \operatorname{Im}(s) \leq \rho_n\} \cap \Gamma$ with appropriate orientations and $\Omega_{\rho_n} = \{s \mid -\rho_n < \operatorname{Im}(s) < \rho_n\} \cap \Omega$.

Using 6 we obtain (for fixed ε)

$$\left| \int_{L_n} \zeta_{\mathcal{T}}(s, \varepsilon) ds \right| \leq \int_{L_n} |\zeta_{\mathcal{T}}(s, \varepsilon)| |ds| \leq K \int_{L_n} \frac{K}{|s|^2} |ds| \leq \frac{K}{\rho_n^2} (c - \sigma_L) \rightarrow 0,$$

as $n \rightarrow \infty$, since the length of L_n is at most $(c - \sigma_L)$ and $\rho_n \rightarrow \infty$ as $n \rightarrow \infty$. Similarly, $\lim_{n \rightarrow \infty} \int_{L'_n} \zeta_{\mathcal{T}}(s, \varepsilon) ds = 0$.

It will be shown in the next lemma that the integral of $\zeta_{\mathcal{T}}$ over Γ absolutely converges, so letting $n \rightarrow \infty$ in 10 gives the desired result. ■

Proof of Lemma 5. Let $R_1 < 0 < R_2$ and let $\Gamma_{R_1}^{R_2}$ be the part of Γ that lies in the strip $\{s \mid R_1 \leq \text{Im}(s) \leq R_2\}$. We will show that, for some $K > 0$ (independent of R_1 and R_2),

$$\int_{\Gamma_{R_1}^{R_2}} |\zeta_{\mathcal{T}}(s, \varepsilon)| |ds| \leq K \varepsilon^{d-\sigma_R}, \quad \text{for } \varepsilon < 1.$$

Suppose $\Gamma_{R_1}^{R_2}$ have non-empty intersection with Γ_k for $k_0 \leq k \leq k_1$. Then

$$\int_{\Gamma_{R_1}^{R_2}} |\zeta_{\mathcal{T}}(s, \varepsilon)| |ds| \leq \sum_{k=k_0}^{k_1} \int_{\Gamma_k} |\zeta_{\mathcal{T}}(s, \varepsilon)| |ds| = \sum_{k=k_0}^{k_1} \sum_{l=1}^4 \int_{\gamma_k^l} |\zeta_{\mathcal{T}}(s, \varepsilon)| |ds|.$$

Recall that, $|\zeta(s)| \leq K$ for $s \in \Gamma$ (by Lemma 3) and by 6,

$$\left| \sum_{m=0}^d \kappa_m \frac{g^{s-m}}{s-m} \right| \leq \frac{K}{|s|^2} \leq \frac{K}{\sigma_L^2 + |\text{Im}s|^2} \quad \text{for } \sigma_L \leq \text{Re}(s) \leq \sigma_R.$$

Now, for $s \in \gamma_k^1$, we have $s = \sigma_R + \mathbf{i}t$, $\beta_{k-1} \leq t \leq \alpha_k$, and

$$\int_{\gamma_k^1} |\zeta_{\mathcal{T}}(s, \varepsilon)| |ds| \leq K \varepsilon^{d-\sigma_R} \int_{\beta_{k-1}}^{\alpha_k} \frac{K}{\sigma_L^2 + t^2} dt.$$

Similarly, $\int_{\gamma_k^3} |\zeta_{\mathcal{T}}(s, \varepsilon)| |ds| \leq K \varepsilon^{d-\sigma_R} \int_{\alpha_k}^{\beta_k} dt / (\sigma_L^2 + t^2)$. Therefore,

$$\sum_{k=k_0}^{k_1} \int_{\gamma_k^1 + \gamma_k^3} |\zeta_{\mathcal{T}}(s, \varepsilon)| |ds| \leq K \varepsilon^{d-\sigma_R} \int_{\beta_{k_0-1}}^{\beta_{k_1}} \frac{dt}{\sigma_L^2 + t^2} \leq K \varepsilon^{d-\sigma_R}. \quad (11)$$

By Lemma 3, $\alpha_{k+1} - \alpha_k \geq \frac{\pi}{\log r_j \Gamma} =: C$, $\alpha_0 < 0$ and $\alpha_1 > 0$. This implies $|\alpha_k| \geq (|k| - 1)C$ and $|\beta_k| \geq (|k| - 1)C$ for all $|k| \geq 2$.

Now for $s \in \gamma_k^2$, we have $s = \sigma + \mathbf{i}\alpha_k$, $\sigma_L \leq \sigma \leq \sigma_R$ and so for $|k| \geq 2$,

$$\int_{\gamma_k^2} |\zeta_{\mathcal{T}}(s, \varepsilon)| |ds| \leq K \varepsilon^{d-\sigma_R} \int_{\gamma_k^2} \frac{|ds|}{|s|^2} \leq K \varepsilon^{d-\sigma_R} \frac{K}{|\alpha_k|^2} (\sigma_R - \sigma_L) \leq \frac{K \varepsilon^{d-\sigma_R}}{(|k| - 1)^2 C^2}.$$

A similar inequality holds when γ_k^2 is replaced with γ_k^4 . Hence

$$\sum_{k=k_0}^{k_1} \int_{\gamma_k^2 + \gamma_k^4} |\zeta_{\mathcal{T}}(s, \varepsilon)| |ds| \leq K \varepsilon^{d-\sigma_R} \quad (12)$$

Combining (11) and (12) we get the desired result. ■

Proof of Lemma 6. Let

$$h(\varepsilon) := \varepsilon^{D-d} \sum_{\omega \in \Omega \cap \mathfrak{D}} \text{res}(\zeta_{\mathcal{T}}(s, \varepsilon); \omega) = \sum_{\omega \in \Omega \cap \mathfrak{D}} \varepsilon^{D-\omega} \text{res}(\zeta(s); \omega) \sum_{m=0}^d \kappa_m \frac{g^{\omega-m}}{\omega-m}.$$

By Lemma 2 and 6

$$|h(\varepsilon)| \leq \sum_{\omega \in \Omega \cap \mathfrak{D}} \frac{1}{|\log r_1|} |\varepsilon^{D-\text{Re}(\omega)}| \frac{K}{|\omega|^2} = K \sum_{\omega \in \Omega \cap \mathfrak{D}} \frac{\varepsilon^{D-\text{Re}(\omega)}}{|\omega|^2}. \quad (13)$$

Let $\Omega_n = \{s \mid -n < \text{Im}(s) < n\} \cap \Omega$, $n = 1, 2, \dots$. By [5, Theorem 3.6], there exists $C, M > 0$ such that the number of poles of $\zeta(s)$ in the strip satisfies

$$Cn - M \leq \#(\mathfrak{D} \cap \{-n < \text{Im}(s) < n\}) \leq Cn + M$$

for all $n = 1, 2, \dots$. Let $\Pi_1 = \Omega_1$ and $\Pi_n = \Omega_n \setminus \Omega_{n-1}$ for $n \geq 2$. Then

$$\#(\mathfrak{D} \cap \Pi_n) = \#(\mathfrak{D} \cap \Omega_n) \setminus (\mathfrak{D} \cap \Omega_{n-1}) \leq Cn + M - (C(n-1) - M) = C + 2M$$

for $n \geq 2$. For $n = 1$, the above inequality clearly holds. Let $\eta > 0$ be arbitrary. We will show that $|h(\varepsilon)| < \eta$ for sufficiently small ε : By 13, for any $n \geq 1$,

$$|h(\varepsilon)| \leq K \sum_{\omega \in \Omega_n \cap \mathfrak{D}} \frac{\varepsilon^{D-\text{Re}(\omega)}}{|\omega|^2} + K \sum_{\omega \in (\Omega \setminus \Omega_n) \cap \mathfrak{D}} \frac{\varepsilon^{D-\text{Re}(\omega)}}{|\omega|^2}. \quad (14)$$

Now, for $\varepsilon < 1$

$$\sum_{\omega \in (\Omega \setminus \Omega_n) \cap \mathfrak{D}} \frac{\varepsilon^{D-\text{Re}(\omega)}}{|\omega|^2} = \sum_{k=n+1}^{\infty} \sum_{\omega \in \Pi_k \cap \mathfrak{D}} \frac{\varepsilon^{D-\text{Re}(\omega)}}{|\omega|^2} \leq \sum_{k=n+1}^{\infty} \frac{C + 2M}{(k-1)^2},$$

since $\varepsilon^{D-\text{Re}(\omega)} < 1$ and for $\omega \in \Pi_k$, ($k \geq 2$), we have $|\omega|^2 \geq |\text{Im}(\omega)|^2 \geq (k-1)^2$. Because of the convergence of the series $\sum_{k=2}^{\infty} \frac{1}{(k-1)^2}$, there exists n_0 such that the second term on the right-hand side of 14 is less than $\eta/2$ for $n = n_0$. To deal with the first term, note that the set $\Omega_{n_0} \cap \mathfrak{D}$ has finitely many elements. Let $\delta := \min\{D - \text{Re}(\omega) \mid \omega \in \Omega_{n_0} \cap \mathfrak{D}\}$. Recall that all the poles of ζ , except the one at $s = D$, have real part less than D (see [5, Theorem 2.17]). Therefore $\delta > 0$ and

$$K \sum_{\omega \in \Omega_{n_0} \cap \mathfrak{D}} \frac{\varepsilon^{D-\text{Re}(\omega)}}{|\omega|^2} \leq K\varepsilon^\delta \sum_{\omega \in \Omega_{n_0} \cap \mathfrak{D}} \frac{1}{|\omega|^2} \leq K\varepsilon^\delta$$

which is less than $\eta/2$ when ε is sufficiently small. ■

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